

Then

$$[B]^T[D][B] = \frac{1}{64} \begin{bmatrix} -(2-y) & 0 & -(1-x) \\ 0 & -(1-x) & -(2-y) \\ \vdots & \vdots & \vdots \\ \vdots & \vdots & \vdots \\ \vdots & \vdots & \vdots \\ \vdots & \vdots & \vdots \\ \vdots & \vdots & \vdots \end{bmatrix} \begin{bmatrix} -c(2-y) & -d(1-x) & \dots & \dots & \dots \\ -d(2-y) & -c(1-x) & \dots & \dots & \dots \\ -e(1-x) & -e(2-y) & \dots & \dots & \dots \end{bmatrix}$$

Therefore

$$K_{11} = \frac{t}{64} \int_{-2}^2 \int_{-1}^1 [c(2-y)^2 + e(1-x)^2] dx dy$$

which gives $K_{11} = \frac{t}{6}(4c + e)$

$$K_{12} = \frac{t}{64} \int_{-2}^2 \int_{-1}^1 [d(2-y)(1-x) + e(1-x)(2-y)] dx dy$$

which gives $K_{12} = \frac{t}{4}(d + e)$.

Solutions to Chapter 7 Problems

S.7.1

Substituting for $((1/\rho_x) + (v/\rho_y))$ and $((1/\rho_y) + (v/\rho_x))$ from Eqs (7.5) and (7.6), respectively in Eqs (7.3)

$$\sigma_x = \frac{Ez}{1-\nu^2} \frac{M_x}{D} \quad \text{and} \quad \sigma_y = \frac{Ez}{1-\nu^2} \frac{M_y}{D} \tag{i}$$

Hence, since, from Eq. (7.4), $D = Et^3/12(1-\nu^2)$, Eqs (i) become

$$\sigma_x = \frac{12zM_x}{t^3} \quad \sigma_y = \frac{12zM_y}{t^3} \tag{ii}$$

The maximum values of σ_x and σ_y will occur when $z = \pm t/2$. Hence

$$\sigma_x(\max) = \pm \frac{6M_x}{t^2} \quad \sigma_y(\max) = \pm \frac{6M_y}{t^2} \tag{iii}$$

Then

$$\sigma_x(\max) = \pm \frac{6 \times 10 \times 10^3}{10^2} = \pm 600 \text{ N/mm}^2$$

$$\sigma_y(\max) = \pm \frac{6 \times 5 \times 10^3}{10^2} = \pm 300 \text{ N/mm}^2$$

S.7.2

From Eq. (7.11) and since $M_{xy} = 0$

$$M_t = \frac{M_x - M_y}{2} \sin 2\alpha \quad (\text{i})$$

M_t will be a maximum when $2\alpha = \pi/2$, i.e. $\alpha = \pi/4$ (45°). Thus, from Eq. (i)

$$M_t(\max) = \frac{10 - 5}{2} = 2.5 \text{ Nm/mm}$$

S.7.3

The relationship between M_n and M_x , M_y and M_{xy} in Eq. (7.10) and between M_t and M_x , M_y and M_{xy} in Eq. (7.11) are identical in form to the stress relationships in Eqs (1.8) and (1.9). Therefore, by deduction from Eqs (1.11) and (1.12)

$$M_I = \frac{M_x + M_y}{2} + \frac{1}{2} \sqrt{(M_x - M_y)^2 + 4M_{xy}^2} \quad (\text{i})$$

and

$$M_{II} = \frac{M_x + M_y}{2} - \frac{1}{2} \sqrt{(M_x - M_y)^2 + 4M_{xy}^2} \quad (\text{ii})$$

Further, Eq. (7.11) gives the inclination of the planes on which the principal moments occur, i.e. when $M_t = 0$. Thus

$$\tan 2\alpha = -\frac{2M_{xy}}{M_x - M_y} \quad (\text{iii})$$

Substituting the values $M_x = 10 \text{ Nm/mm}$, $M_y = 5 \text{ Nm/mm}$ and $M_{xy} = 5 \text{ Nm/mm}$ in Eqs (i), (ii) and (iii) gives

$$M_I = 13.1 \text{ Nm/mm}$$

$$M_{II} = 1.9 \text{ Nm/mm}$$

and

$$\alpha = -31.7^\circ \quad \text{or} \quad 58.3^\circ$$

The corresponding principal stresses are obtained directly from Eqs (iii) of S.7.1. Hence

$$\sigma_I = \pm \frac{6 \times 13.1 \times 10^3}{10^2} = \pm 786 \text{ N/mm}^2$$

$$\sigma_{II} = \pm \frac{6 \times 1.9 \times 10^3}{10^2} = \pm 114 \text{ N/mm}^2$$

S.7.4

From the deflection equation

$$\frac{\partial^2 w}{\partial x^2} = -\frac{q_0 a^2}{D\pi^2} \left(1 + A \cosh \frac{\pi y}{a} + B \frac{\pi y}{a} \sinh \frac{\pi y}{a} \right) \sin \frac{\pi x}{a}$$

$$\frac{\partial^2 w}{\partial y^2} = \frac{q_0 a^2}{D\pi^2} \left(A \cosh \frac{\pi y}{a} + 2B \cosh \frac{\pi y}{a} + B \frac{\pi y}{a} \sinh \frac{\pi y}{a} \right) \sin \frac{\pi x}{a}$$

Now $w = 0$ and $M_x = 0$ at $x = 0$ and a . From Eq. (7.7) this is satisfied implicitly.

Also $w = 0$ and $M_y = 0$ at $y = \pm a$ so that, from the deflection equation

$$0 = \frac{q_0 a^4}{D\pi^4} (1 + A \cosh \pi + B\pi \sinh \pi) \sin \frac{\pi x}{a}$$

i.e.

$$1 + A \cosh \pi + B\pi \sinh \pi = 0 \quad (i)$$

Also, from Eq. (7.8)

$$0 = -\frac{q_0 a^2}{D\pi^2} [(A \cosh \pi + 2B \cosh \pi + B\pi \sinh \pi) - 0.3(1 + A \cosh \pi + B\pi \sinh \pi)] \sin \frac{\pi x}{a}$$

or

$$0 = -0.3 + 0.7A \cosh \pi + 2B \cosh \pi + 0.7B\pi \sinh \pi \quad (ii)$$

Solving Eqs (i) and (ii)

$$A = -0.2213 \quad B = 0.0431$$

S.7.5

The deflection is zero at $x = a/2$, $y = a/2$. Then, from the deflection equation

$$0 = \frac{a^4}{4} - \frac{3}{2}a^4(1 - \nu) - \frac{3}{4}a^4\nu + A$$

Hence

$$A = \frac{a^4}{4}(5 - 3\nu)$$

The central deflection, i.e. at $x = 0, y = 0$ is then

$$\begin{aligned} &= \frac{qa^4}{96(1-\nu)D} \times \frac{1}{4}(5 - 3\nu) \\ &= \frac{qa^4}{384D} \left(\frac{5 - 3\nu}{1 - \nu} \right) \end{aligned}$$

S.7.6

From the equation for deflection

$$\begin{aligned} \frac{\partial^4 w}{\partial x^4} &= w_0 \left(\frac{\pi}{a} \right)^4 \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a} \\ \frac{\partial^4 w}{\partial y^4} &= w_0 \left(\frac{3\pi}{a} \right)^4 \cos \frac{\pi x}{a} \sin \frac{3\pi y}{a} \\ \frac{\partial^4 w}{\partial x^2 \partial y^2} &= w_0 \left(\frac{\pi}{a} \right)^2 \left(\frac{3\pi}{a} \right)^2 \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a} \end{aligned}$$

Substituting in Eq. (7.20)

$$\frac{q(x, y)}{D} = w_0 \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a} (1 + 2 \times 9 + 81) \left(\frac{\pi}{a} \right)^4$$

i.e.

$$q(x, y) = w_0 D 100 \frac{\pi^4}{a^4} \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a}$$

From the deflection equation

$$w = 0 \quad \text{at } x = \pm a/2, y = \pm a/2$$

The plate is therefore supported on all four edges.

Also

$$\begin{aligned} \frac{\partial w}{\partial x} &= -w_0 \frac{\pi}{a} \sin \frac{\pi x}{a} \cos \frac{3\pi y}{a} \\ \frac{\partial w}{\partial y} &= -w_0 \frac{3\pi}{a} \cos \frac{\pi x}{a} \sin \frac{3\pi y}{a} \end{aligned}$$

When

$$x = \pm \frac{a}{2} \quad \frac{\partial w}{\partial x} \neq 0$$

$$y = \pm \frac{a}{2} \quad \frac{\partial w}{\partial y} \neq 0$$

The plate is therefore not clamped on its edges.

Further

$$\frac{\partial^2 w}{\partial x^2} = -w_0 \left(\frac{\pi}{a}\right)^2 \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a}$$

$$\frac{\partial^2 w}{\partial y^2} = -w_0 \left(\frac{3\pi}{a}\right)^2 \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a}$$

Substituting in Eq. (7.7)

$$M_x = -Dw_0 \left(\frac{\pi}{a}\right)^2 \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a} (-1 - 9\nu) \quad (\text{i})$$

Similarly, from Eq. (7.8)

$$M_y = w_0 D \left(\frac{\pi}{a}\right)^2 \cos \frac{\pi x}{a} \cos \frac{3\pi y}{a} (9 + \nu) \quad (\text{ii})$$

Then, at $x = \pm a/2$, $M_x = 0$ (from Eq. (i)) and at $y = \pm a/2$, $M_y = 0$ (from Eq. (ii)).

The plate is therefore simply supported on all edges.

The corner reactions are given by

$$2D(1 - \nu) \frac{\partial^2 w}{\partial x \partial y} \quad (\text{see Eq. (7.14)})$$

Then, since

$$\frac{\partial^2 w}{\partial x \partial y} = w_0 \frac{\pi}{a} \frac{3\pi}{a} \sin \frac{\pi x}{a} \sin \frac{3\pi y}{a} \quad \text{at } x = a/2, y = a/2$$

$$\text{Corner reactions} = -6w_0 D \left(\frac{\pi}{a}\right)^2 (1 - \nu)$$

From Eqs (7.7) and (7.8) and the above, at the centre of the plate

$$M_x = w_0 D \left(\frac{\pi}{a}\right)^2 (1 + 9\nu), \quad M_y = w_0 D \left(\frac{\pi}{a}\right)^2 (9 + \nu).$$

S.7.7

Substituting $q(x, y) = q_0 x/a$ in Eq. (7.29) and noting that the plate is square and of side a

$$a_{mn} = \frac{4}{a^2} \int_0^a \int_0^a q_0 \frac{x}{a} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{a} dx dy$$

i.e.

$$a_{mn} = \frac{4q_0}{a^3} \int_0^a x \sin \frac{m\pi x}{a} \left[-\frac{a}{n\pi} \cos \frac{n\pi y}{a} \right]_0^a dx$$

Hence

$$a_{mn} = -\frac{4q_0}{a^2 n\pi} \int_0^a x \sin \frac{m\pi x}{a} (\cos n\pi - 1) dx$$

The term in brackets is zero when n is even and equal to -2 when n is odd. Thus

$$a_{mn} = \frac{8q_0}{a^2 n\pi} \int_0^a x \sin \frac{m\pi x}{a} dx \quad (n \text{ odd}) \quad (i)$$

Integrating Eq. (i) by parts

$$a_{mn} = \frac{8q_0}{a^2 n\pi} \left[-x \frac{a}{m\pi} \cos \frac{m\pi x}{a} + \int \frac{a}{m\pi} \cos \frac{m\pi x}{a} dx \right]_0^a$$

i.e.

$$a_{mn} = \frac{8q_0}{am n\pi^2} \left[-x \cos \frac{m\pi x}{a} + \frac{a}{m\pi} \sin \frac{m\pi x}{a} \right]_0^a$$

The second term in square brackets is zero for all integer values of m . Thus

$$a_{mn} = \frac{8q_0}{am n\pi^2} (-a \cos m\pi)$$

The term in brackets is positive when m is odd and negative when m is even. Thus

$$a_{mn} = \frac{8q_0}{m n\pi^2} (-1)^{m+1}$$

Substituting for a_{mn} in Eq. (7.30) gives the displaced shape of the plate, i.e.

$$w = \frac{1}{\pi^4 D} \sum_{m=1,2,3}^{\infty} \sum_{n=1,3,5}^{\infty} \frac{8q_0(-1)^{m+1}}{m n\pi^2 \left[\left(\frac{m^2}{a^2} \right) + \left(\frac{n^2}{a^2} \right) \right]^2} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{a}$$

or

$$w = \frac{8q_0 a^4}{\pi^6 D} \sum_{m=1,2,3}^{\infty} \sum_{n=1,3,5}^{\infty} \frac{(-1)^{m+1}}{m n(m^2 + n^2)^2} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{a}$$

S.7.8

The boundary conditions which must be satisfied by the equation for the displaced shape of the plate are $w = 0$ and $\partial w/\partial n = 0$ at all points on the boundary; n is a direction normal to the boundary at any point.

The equation of the ellipse representing the boundary is

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1 \quad (\text{i})$$

Substituting for $x^2/a^2 + y^2/b^2$ in the equation for the displaced shape clearly gives $w = 0$ for all values of x and y on the boundary of the plate. Also

$$\frac{\partial w}{\partial n} = \frac{\partial w}{\partial x} \frac{\partial x}{\partial n} + \frac{\partial w}{\partial y} \frac{\partial y}{\partial n} \quad (\text{ii})$$

Now

$$w = w_0 \left(1 - \frac{x^2}{a^2} - \frac{y^2}{b^2} \right)^2$$

so that

$$\frac{\partial w}{\partial x} = -\frac{4w_0x}{a^2} \left(1 - \frac{x^2}{a^2} - \frac{y^2}{b^2} \right) \quad (\text{iii})$$

and

$$\frac{\partial w}{\partial y} = -\frac{4w_0y}{b^2} \left(1 - \frac{x^2}{a^2} - \frac{y^2}{b^2} \right) \quad (\text{iv})$$

From Eqs (i), (ii) and (iv) it can be seen that $\partial w/\partial x$ and $\partial w/\partial y$ are zero for all values of x and y on the boundary of the plate. It follows from Eq. (ii) that $\partial w/\partial n = 0$ at all points on the boundary of the plate. Thus the equation for the displaced shape satisfies the boundary conditions.

From Eqs (iii) and (iv)

$$\frac{\partial^4 w}{\partial x^4} = \frac{24w_0}{a^4} \quad \frac{\partial^4 w}{\partial y^4} = \frac{24w_0}{b^4} \quad \frac{\partial^4 w}{\partial x^2 \partial y^2} = \frac{8w_0}{a^2 b^2}$$

Substituting these values in Eq. (7.20)

$$w_0 \left(\frac{24}{a^4} + \frac{16}{a^2 b^2} + \frac{24}{b^4} \right) = \frac{p}{D}$$

whence

$$w_0 = \frac{p}{8D \left(\frac{3}{a^4} + \frac{2}{a^2 b^2} + \frac{3}{b^4} \right)}$$

Now substituting for D from Eq. (7.4)

$$w_0 = \frac{3p(1 - \nu^2)}{2Et^3 \left(\frac{3}{a^4} + \frac{2}{a^2b^2} + \frac{3}{b^4} \right)} \quad (\text{v})$$

From Eqs (7.3), (7.5) and (7.7)

$$\sigma_x = -\frac{Ez}{1 - \nu^2} \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right) \quad (\text{vi})$$

and from Eqs (7.3), (7.6) and (7.8)

$$\sigma_y = -\frac{Ez}{1 - \nu^2} \left(\frac{\partial^2 w}{\partial y^2} + \nu \frac{\partial^2 w}{\partial x^2} \right) \quad (\text{vii})$$

From Eqs (iii) and (iv)

$$\frac{\partial^2 w}{\partial x^2} = -\frac{4w_0}{a^2} \left(1 - \frac{3x^2}{a^2} - \frac{y^2}{b^2} \right) \quad \frac{\partial^2 w}{\partial y^2} = -\frac{4w_0}{b^2} \left(1 - \frac{x^2}{a^2} - \frac{3y^2}{b^2} \right)$$

Substituting these expressions in Eq. (vi) and noting that the maximum values of direct stress occur at $z = \pm t/2$

$$\sigma_x(\text{max}) = \pm \frac{Et}{2(1 - \nu^2)} \left[-\frac{4w_0}{a^2} \left(1 - \frac{3x^2}{a^2} - \frac{y^2}{b^2} \right) - \frac{4w_0\nu}{b^2} \left(1 - \frac{x^2}{a^2} - \frac{3y^2}{b^2} \right) \right] \quad (\text{viii})$$

At the centre of the plate, $x = y = 0$. Then

$$\sigma_x(\text{max}) = \pm \frac{2Et w_0}{(1 - \nu^2)} \left(\frac{1}{a^2} + \frac{\nu}{b^2} \right) \quad (\text{ix})$$

Substituting for w_0 in Eq. (ix) from Eq. (v) gives

$$\sigma_x(\text{max}) = \pm \frac{3pa^2b^2(b^2 + \nu a^2)}{t^2(3b^4 + 2a^2b^2 + 3a^4)} \quad (\text{x})$$

Similarly

$$\sigma_y(\text{max}) = \pm \frac{3pa^2b^2(a^2 + \nu b^2)}{t^2(3b^4 + 2a^2b^2 + 3a^4)} \quad (\text{xi})$$

At the ends of the minor axis, $x = 0$, $y = b$. Thus, from Eq. (viii)

$$\sigma_x(\text{max}) = \pm \frac{2Et w_0}{(1 - \nu^2)} \left(\frac{1}{a^2} - \frac{1}{a^2} + \frac{\nu}{b^2} - \frac{3\nu}{b^2} \right)$$

i.e.

$$\sigma_x(\text{max}) = \pm \frac{4Et w_0 \nu}{b^2(1 - \nu^2)} \quad (\text{xii})$$

Again substituting for w_0 from Eq. (v) in Eq. (xii)

$$\sigma_x(\max) = \pm \frac{6pa^4b^2}{t^2(3b^4 + 2a^2b^2 + 3a^4)}$$

Similarly

$$\sigma_y(\max) = \pm \frac{6pb^4a^2}{t^2(3b^4 + 2a^2b^2 + 3a^4)}$$

S.7.9

The potential energy, V , of the load W is given by

$$V = -Ww$$

i.e.

$$V = -W \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b}$$

Therefore, it may be deduced from Eq. (7.47) that the total potential energy, $U + V$, of the plate is

$$U + V = \frac{D}{2} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn}^2 \frac{\pi^4 ab}{4} \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2 - W \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b}$$

From the principle of the stationary value of the total potential energy

$$\frac{\partial(U + V)}{\partial A_{mn}} = DA_{mn} \frac{\pi^4 ab}{4} \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2 - W \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b} = 0$$

Hence

$$A_{mn} = \frac{4W \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b}}{\pi^4 Dab \left[\left(\frac{m^2}{a^2} \right) + \left(\frac{n^2}{b^2} \right) \right]^2}$$

so that the deflected shape is obtained.

S.7.10

From Eq. (7.45) the potential energy of the in-plane load, N_x , is

$$-\frac{1}{2} \int_0^a \int_0^b N_x \left(\frac{\partial w}{\partial x} \right)^2 dx dy$$

The combined potential energy of the in-plane load, N_x , and the load, W , is then, from S.7.9

$$V = -\frac{1}{2} \int_0^a \int_0^b N_x \left(\frac{\partial w}{\partial x} \right)^2 dx dy - W \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b}$$

or, since,

$$\frac{\partial w}{\partial x} = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \frac{m\pi}{a} \cos \frac{m\pi x}{a} \sin \frac{n\pi y}{b}$$

$$\begin{aligned} V = & -\frac{1}{2} \int_0^a \int_0^b N_x \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn}^2 \frac{m^2\pi^2}{a^2} \cos^2 \frac{m\pi x}{a} \sin^2 \frac{n\pi y}{b} dx dy \\ & - W \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b} \end{aligned}$$

i.e.

$$V = -\frac{ab}{8} N_x \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn}^2 \frac{m^2\pi^2}{a^2} - W \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b}$$

Then, from Eq. (7.47), the total potential energy of the plate is

$$\begin{aligned} U + V = & \frac{D}{2} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn}^2 \frac{\pi^4 ab}{4} \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2 - \frac{ab}{8} N_x \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn}^2 \frac{m^2\pi^2}{a^2} \\ & - W \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b} \end{aligned}$$

Then, from the principle of the stationary value of the total potential energy

$$\frac{\partial(U + V)}{\partial A_{mn}} = DA_{mn} \frac{\pi^4 ab}{4} \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2 - \frac{ab}{4} N_x A_{mn} \frac{m^2\pi^2}{a^2} - W \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b} = 0$$

from which

$$A_{mn} = \frac{4W \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b}}{abD\pi^4 \left[\left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2 - \frac{m^2 N_x}{\pi^2 a^2 D} \right]}$$

S.7.11

The guessed form of deflection is

$$w = A_{11} \left(1 - \frac{4x^2}{a^2}\right) \left(1 - \frac{4y^2}{a^2}\right) \quad (i)$$

Clearly when $x = \pm a/2$, $w = 0$ and when $y = \pm a/2$, $w = 0$. Therefore, the equation for the displaced shape satisfies the displacement boundary conditions.

From Eq. (i)

$$\frac{\partial^2 w}{\partial x^2} = -8 \frac{A_{11}}{a^2} \left(1 - \frac{4y^2}{a^2}\right) \quad \frac{\partial^2 w}{\partial y^2} = -8 \frac{A_{11}}{a^2} \left(1 - \frac{4x^2}{a^2}\right)$$

Substituting in Eq. (7.7)

$$M_x = -\frac{8A_{11}D}{a^2} \left[1 - \frac{4y^2}{a^2} + \nu \left(1 - \frac{4x^2}{a^2}\right)\right]$$

Clearly, when $x = \pm a/2$, $M_x \neq 0$ and when $y = \pm a/2$, $M_x \neq 0$. Similarly for M_y . Thus the assumed displaced shape does not satisfy the condition of zero moment at the simply supported edges.

From Eq. (i)

$$\frac{\partial^2 w}{\partial x \partial y} = \frac{64A_{11}xy}{a^4}$$

Substituting for $\partial^2 w/\partial x^2$, $\partial^2 w/\partial y^2$, $\partial^2 w/\partial x \partial y$ and w in Eq. (7.46) and simplifying gives

$$\begin{aligned} U + V = \int_{-a/2}^{a/2} \int_{-a/2}^a & \left\{ \frac{32A_{11}^2 D}{a^4} \left[4 - \frac{16}{a^2}(x^2 + y^2) + \frac{16}{a^4}(x^4 + 2x^2y^2 + y^4) - 1.4 \right. \right. \\ & \left. \left. + \frac{5.6}{a^2}(x^2 + y^2) + \frac{67.2x^2y^2}{a^4} \right] \right. \\ & \left. - q_0 A_{11} \left(1 - \frac{4x^2}{a^2} - \frac{4y^2}{a^2} + \frac{16x^2y^2}{a^4} \right) \right\} dx dy \end{aligned}$$

from which

$$U + V = \frac{62.4A_{11}^2 D}{a^2} - \frac{4q_0 A_{11} a^2}{9}$$

From the principle of the stationary value of the total potential energy

$$\frac{\partial(U + V)}{\partial A_{11}} = \frac{124.8A_{11}D}{a^2} - \frac{4q_0 a^2}{9} = 0$$

Hence, since $D = Et^3/12(1 - \nu^2)$

$$A_{11} = 0.0389q_0 a^4 / Et^3$$

S.7.12

From Eq. (7.36) the deflection of the plate from its initial curved position is

$$w_1 = B_{11} \sin \frac{\pi x}{a} \sin \frac{\pi y}{b}$$

in which

$$B_{11} = \frac{A_{11}N_x}{\frac{\pi^2 D}{a^2} \left(1 + \frac{a^2}{b^2}\right)^2 - N_x}$$

The total deflection, w , of the plate is given by

$$w = w_1 + w_0$$

i.e.

$$w = \left[\frac{A_{11}N_x}{\frac{\pi^2 D}{a^2} \left(1 + \frac{a^2}{b^2}\right)^2 - N_x} + A_{11} \right] \sin \frac{\pi x}{a} \sin \frac{\pi y}{b}$$

i.e.

$$w = \frac{A_{11}}{1 - \frac{N_x a^2}{\pi^2 D} / \left(1 + \frac{a^2}{b^2}\right)^2} \sin \frac{\pi x}{a} \sin \frac{\pi y}{b}$$

Solutions to Chapter 8 Problems

S.8.1

The forces on the bar AB are shown in Fig. S.8.1 where

$$M_B = K \left(\frac{dv}{dz} \right)_B \quad (\text{i})$$

and P is the buckling load.

From Eq. (8.1)

$$EI \frac{d^2 v}{dz^2} = -Pv \quad (\text{ii})$$

The solution of Eq. (ii) is

$$v = A \cos \mu z + B \sin \mu z \quad (\text{iii})$$